Computing separation variables on a computer

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Orthogonal coordinate systems

Different coordinate systems are convenient for different problems. For example, spherical coordinate system is convenient for a problem of a particle in a central force field, while oblate spheroidal system can be convenient for the Jacobi-Calogero inverse square model.





spherical

oblate spheroidal

Using separation variables in numerical simulations

- Reduces computation time
- Easier parallelizing
- Easier control of the correctness through integrals of motion

Theoretical overview

Q - riemannian manifold with local coordinates $q = (q^1, q^2, \dots, q^n)$ and positive definite metric tensor **G**.

Examples:





sphere

ellipsoid

Natural Hamilton function

$$H = T(p,q) + V(q) = \sum_{i,j=1}^{n} g^{ij}(q) p_i p_j + V(q)$$

Separation of variables for the Hamilton-Jacobi equation

Hamilton-Jacobi equationH(p,q) = EAdditive Action $S(Q, \alpha_1, \dots, \alpha_n) \equiv \sum_{i=1}^n S_i(Q_i, \alpha_1, \dots, \alpha_n)$ Conjugated momenta $P_i = \frac{\partial S_i(Q_i, \alpha_1, \dots, \alpha_n)}{\partial Q_i}$ Separated equations $\Phi_i(Q_i, P_i, \alpha_1, \dots, \alpha_n) = 0$

Levi-Civita criteria

The Hamilton-Jacobi equation admits separation of variables if the Hamilton function H(P,Q) satisfies the following n(n-1)/2 equations with $i \neq j$

 $\frac{\partial H}{\partial P_{j}} \frac{\partial H}{\partial P_{i}} \frac{\partial^{2} H}{\partial Q_{j} \partial Q_{i}} - \frac{\partial H}{\partial P_{j}} \frac{\partial H}{\partial Q_{i}} \frac{\partial^{2} H}{\partial Q_{j} \partial P_{i}}$ $- \frac{\partial H}{\partial Q_{j}} \frac{\partial H}{\partial P_{i}} \frac{\partial^{2} H}{\partial P_{j} \partial Q_{i}} + \frac{\partial H}{\partial Q_{j}} \frac{\partial H}{\partial Q_{i}} \frac{\partial^{2} H}{\partial P_{j} \partial P_{i}} = 0.$

Separating variables with coordinate transformations

Point transformations:

$$Q = f(q),$$
 $P = f'(q) p,$ $f = (f_1, ..., f_n)$

The Hamiltonian is covariant w.r.t point transformations

$$H = T(P,Q) + V(Q) = \sum_{i,j=1}^{n} g^{ij}(Q) P_i P_j + V(Q).$$

Killing Tensors

 \mathcal{Q} - riemannian manifold with a metric tensor $\mathbf{G} = g^{ij}\partial_i \otimes \partial_j, \qquad \partial_k \equiv \frac{\partial}{\partial q^k}$

The correspondence between tensor \mathbf{K} and polynomial P_K :

$$\mathbf{K} = (K^{i \dots j}) \qquad \longleftrightarrow \qquad P_K = K^{i \dots j} p_i \cdots p_j.$$

Definition 1 Killing Tensor K of rank ℓ is a symmetric $(\ell, 0)$ tensor in the space Q satisfying tensor Killing equation

$$[\mathbf{K},\mathbf{G}] = \mathbf{0} \qquad \Longleftrightarrow \qquad \{P_K,P_G\} = \mathbf{0}.$$

Definition 2 Conformal Killing Tensor L of rank ℓ is a symmetric $(\ell, 0)$ tensor in the space Q, satisfying tensor equation

$$\{P_L, P_G\} = cP_G.$$

Definition 3 Conformal tensor L with simple eigenvalues and zero Nijenhuis torsion

$$T_{ij}^m \equiv 2L_i^\alpha \partial_\alpha L_j^m - 2L_\alpha^m \partial_i L_j^\alpha = 0.$$

is called L-tensor, or Benenti tensor.

Separating variables for L-systems

Theorem 1 [Benenti] Eigenvalues Q_i of tensor L are separation variables for the integrable system.

Integrable systems admitting separation in such variables are called L-systems, or Benenti systems.

For given L we can define a recursion operator N, i.e. canonical lifting of tensor L on a cotangent bundle T^*Q according to the rule

$$\mathbf{N} \frac{\partial}{\partial q^k} = \sum_{i=1}^n L_k^i \frac{\partial}{\partial q^i} + \sum_{ij} p_j \left(\frac{\partial L_i^j}{\partial q^k} - \frac{\partial L_k^j}{\partial q^i} \right) \frac{\partial}{\partial p_i},$$
$$\mathbf{N} \frac{\partial}{\partial p_k} = \sum_{i=1}^n L_i^k \frac{\partial}{\partial p_i}.$$

Recurrent equations for the integrals of motion H_m :

$$dH_{m+1} = N^* dH_m + \sigma_{m+1} dH, \qquad m = 1, ..., n-1, \qquad H_n \equiv 0.$$

Here σ_m are the coefficients of the characteristic polynomial of L: det($\lambda I - L$) = $\sum_{m=0}^{n} \sigma_m \lambda^{n-m}$

Algorithm of calculating separation variables

1. Build and solve equations for L

$$d(i_{X_T}d\theta - Td\sigma_1) = 0$$

$$d(i_{X_V}d\theta - Vd\sigma_1) = 0$$

Here i is a hook operator, $\sigma_1 = \operatorname{tr} \mathbf{L}$ and $\theta = \sum_{i,j=1}^n L_j^i p_i dq^j$ is an
L-deformation of standard Liouville 1-form $\theta_0 = \sum p_j dq^j$.

- 2. Calculate eigenvalues of L separation variables
- 3. Solve a chain of equations $dH_{m+1} = \mathbf{N}^* dH_m + \sigma_{m+1} dH$ for the integrals of motion.

Implementation (Maple worksheet)

Applying the programme to some integrable systems 1. Neumann system The Neumann system describes the motion of a particle on a sphere under the influence of a

quadratic potential $V(x) = \frac{1}{2} \sum_{i=1}^{N+1} \alpha_i x_i^2$, where $\alpha_i \neq \alpha_j$

Spheroconical coordinates:
$$\sum_{i=1}^{N+1} \frac{x_i^2}{\lambda - \alpha_i} =$$

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Manual transformation of variables

$$X_i = x_i^2, i = 1, ..., N$$

$$T = 2\sum_{a} X_a (1 - X_a) Y_a^2 - 4 \sum_{a < b} X_a X_b Y_a Y_b$$

$$V = \frac{1}{2} \sum_{a} (\alpha_a - \alpha_{N+1}) X_a$$

2. Holt system

$$H(x) = \frac{1}{2}(p_X^2 + p_Y^2) + aX^{-\frac{2}{3}}(\frac{3b}{4}X^2 + Y^2 + c)$$

Manual transformation of variables

$$a \to 4(\frac{3}{2})^{1/3}a, \quad c \to \frac{c}{3a}$$

$$X = \frac{2}{3}x^{3/2}, \quad p_X = p_x/\sqrt{x}$$

$$Y = -\frac{1}{2\sqrt{3a}}p_y \quad p_Y = 2\sqrt{3a}y$$

$$H = \frac{p_x^2 + p_y^2}{2x} + 2a(bx^2 + 3y^2) + \frac{2c}{x}$$

Ways of improving the performance of the programme

• Watching after and modifying the intermediate results of the programme

• Using special PDE solvers for overdetermined systems

Ways of broadening the class of systems separable by the programme

• Performing manual transformations of variables

 \bullet Solving the systems for ${\bf N}$ instead of those for ${\bf L}$

 Adding typical non-point transformations of variables with arbitrary parameters to the equations for L